Reduction Groups and Darboux Transformations

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- Lax Operators
- Reduction Groups
- Darboux Transformations
- Integrable partial difference equations

What is a Lax operator?

• For my purposes I will consider a pair of Lax operators to be differential operators of the form:

$$L = \partial_x - X(\lambda, x, t)$$
$$S = \partial_t - T(\lambda, x, t)$$

where X and T are matrix functions of x ,t and λ , here λ is called the spectral parameter.

- Lax pairs are important in the study of non-linear pdes.
- Suppose we have a non-zero vector Ψ such that

$$L\Psi = \partial_x \Psi - X\Psi = 0$$

$$S\Psi = \partial_t \Psi - T\Psi = 0.$$

• This system is consistent iff $\partial_x \partial_t \Psi = \partial_t \partial_x \Psi$ which is iff

$$X_t - T_x + [X, T] = 0$$
 (Zero Curvature Condition).

Example

For Lax operators having X,T as

$$X = \left(\begin{array}{cc} -i\lambda & u \\ -1 & i\lambda \end{array}\right)$$

$$T = \begin{pmatrix} -4i\lambda^3 + 2iu\lambda - u_x & 4u\lambda^2 + 2iu_x\lambda - 2u^2 - u_{xx} \\ -4\lambda^2 + 2u & 4i\lambda^3 - 2iu\lambda + u_x \end{pmatrix},$$

the Zero Curvature condition becomes:

$$X_t - T_x + [X, T] = \begin{pmatrix} 0 & u_t + 6uu_x + u_{xxx} \\ 0 & 0 \end{pmatrix} = 0$$

so holds iff u satisfies the KdV equation.

AKNS Scheme

In fact the previous example is a specialisation of the more general Lax pair with:

$$X = \begin{pmatrix} -i\lambda & u \\ v & i\lambda \end{pmatrix}$$
$$T = \begin{pmatrix} A & B \\ C & -A \end{pmatrix}.$$

- Take $A = \sum_{i=0}^{n} a_i \lambda^i$, $B = \sum_{i=0}^{n} b_i \lambda^i$, $C = \sum_{i=0}^{n} c_i \lambda^i$.
- Set the coefficients of λ to zero in the Zero Curvature condition
- Solve the resulting system to produce Integrable non-linear pdes for different integers *n*.
- For n = 2 we can find the non-linear Schrodinger equation.
- For n = 3 we may recover the KdV equation.
- For n = -1 the Sin-Gordan equation.
- This scheme is named after Ablowitz, Kaup, Newell and Segur.

Reduction Problem

In theory we could consider Lax operators of the form:

$$L = \partial_x + X_0 + \sum_{i=1}^n \frac{1}{(\lambda - \gamma_i)^{\alpha_i}} X_i$$

 $S = \partial_t + T_0 + \sum_{i=1}^n \frac{1}{(\lambda - \mu_i)^{\beta_i}} T_i$

where the entries of the X_i , T_i are functions of x and t.

- Too general to be useful.
- We need a way to reduce the generality of X and T.
- We do this by assuming that X_i and T_i are invariant with respect to the action of a finite group of transformations, the Reduction Group.

Automorphic Lie Algebras

X, T ∈ 𝔅_λ(Γ) = R_λ(Γ) ⊗_C 𝔅 where R_λ(Γ) is the ring of rational functions in λ with poles in Γ, 𝔅 is a finite dimensional semi simple Lie Algebra.

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$$\Gamma = \{\infty\}$$
 gives that $R_{\lambda}(\Gamma) = \mathbb{C}[\lambda]$

- $\Gamma = \{0,\infty\}$ gives that $R_{\lambda}(\Gamma) = \mathbb{C}[\lambda,\lambda^{-1}]$
- Letting $G \subset Aut(\mathfrak{U}_{\lambda}(\Gamma))$ consider the set of elements $\mathfrak{U}_{\lambda}^{G} = \{U \in \mathfrak{U}_{\lambda}(\Gamma) : g(U) = U, \forall g \in G\}.$
- We call this subalgebra the Automorphic Lie algebra corresponding to the group G and the set Γ. G is called the reduction group.
- This terminology was introduced by S. Lombardo and A.V. Mikhailov in "Reduction Groups and Automorphic Lie Algebras" Commun. Math. Phys. 258, 179-202 (2005).

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Types of reduction group

- We consider simultaneous automorphisms of $R_{\lambda}(\Gamma)$ and \mathfrak{U} .
- Automorphisms of $R_{\lambda}(\Gamma)$ are fractional linear transformations. Finite subgroups of the group of fractional linear transformations have been classified by Felix Klein. The complete list is given by:
 - \mathbb{Z}_N Cyclic groups
 - \mathbb{D}_N Dihedral groups
 - \mathbb{T} Tetrahedral group
 - \mathbb{O} Octahedral group
 - $\mathbb I$ lcosahedral group.
- For $\mathfrak{U} = \mathfrak{sl}(N, \mathbb{C})$ automorphisms consist of:
 - Inner automorphisms $a \mapsto QaQ^{-1}$ for N = 2
 - Inner and outer automorphisms $a \mapsto -a^{tr}$ for $N \ge 3$.
- We find irreducible projective representations. These consist of 2x2 representations for all reduction groups, 3x3 representations for T, O and I, 4x4 representations for O and I, 5x5 and 6x6 for I.

Structure of Automorphic Lie algebras

- We will consider Γ = G(γ), so a single orbit. This means we have two situations, Γ is either a generic orbit or a degenerated orbit.
- Automorphic Lie algebras are constructed by taking the group averages $\langle \frac{\mathbf{e}_i}{(\lambda-\gamma)^{k_i}} \rangle_G = \frac{1}{|G|} \sum_{g \in G} g(\frac{\mathbf{e}_i}{(\lambda-\gamma)^{k_i}})$, where $\mathbf{e}_1, \mathbf{e}_2, \ldots, \mathbf{e}_m$ are a basis of $\mathfrak{U}, \gamma \in \Gamma$ and k_i is chosen such that the sum is non-zero.

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Example

The \mathbb{D}_2 reduction group can be defined by the generators:

$$g_s: \mathbf{a}(\lambda) \to \mathbf{Q}_s \mathbf{a}(\sigma_s^{-1}(\lambda)) \mathbf{Q}_s^{-1}$$

 $g_r: \mathbf{a}(\lambda) \to \mathbf{Q}_r \mathbf{a}(\sigma_r^{-1}(\lambda)) \mathbf{Q}_r^{-1}$

where

$$\sigma_{s}(\lambda) = -\lambda , \sigma_{r}(\lambda) = 1/\lambda$$
$$\mathbf{Q}_{s} = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}, \mathbf{Q}_{r} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}.$$

In this case the group average is given by

$$< \mathbf{a}(\lambda) >_{\mathbb{D}_2} = rac{1}{4} (\mathbf{a}(\lambda) + \mathbf{Q}_s \mathbf{a}(-\lambda) \mathbf{Q}_s^{-1} + \mathbf{Q}_r \mathbf{a}(-\lambda^{-1}) \mathbf{Q}_r^{-1} + \mathbf{Q}_r \mathbf{Q}_s \mathbf{a}(-\lambda^{-1}) \mathbf{Q}_s^{-1} \mathbf{Q}_r^{-1})$$

Example continued...

Taking the standard basis for $sl(2,\mathbb{C})$

$$\mathbf{e}_1 = \left(\begin{array}{cc} 0 & 1 \\ 0 & 0 \end{array} \right), \mathbf{e}_2 = \left(\begin{array}{cc} 0 & 0 \\ 1 & 0 \end{array} \right), \mathbf{e}_3 = \left(\begin{array}{cc} 1 & 0 \\ 0 & -1 \end{array} \right),$$

we obtain the following:

$$\begin{split} \hat{\mathbf{a}}_1 = & < \frac{\mathbf{e}_1}{\lambda - \gamma} >_{\mathbb{D}_2} = \begin{pmatrix} 0 & \frac{\lambda}{2(\lambda^2 - \gamma^2)} \\ \frac{\lambda}{2(1 - \lambda^2 \gamma^2)} & 0 \end{pmatrix} \\ \hat{\mathbf{a}}_2 = & < \frac{\mathbf{e}_2}{\lambda - \gamma} >_{\mathbb{D}_2} = \begin{pmatrix} 0 & \frac{\lambda}{2(1 - \lambda^2 \gamma^2)} \\ \frac{\lambda}{2(\lambda^2 - \gamma^2)} & 0 \end{pmatrix} \\ \hat{\mathbf{a}}_3 = & < \frac{\mathbf{e}_3}{\lambda - \gamma} >_{\mathbb{D}_2} = \frac{\gamma(1 - \lambda^4)}{2(\lambda^2 - \gamma^2)(1 - \lambda^2 \gamma^2)} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}. \end{split}$$

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Example continued...

Setting

$$\mathbf{a}_1^n = 4\hat{\mathbf{a}}_1 J^n$$
$$\mathbf{a}_2^n = 4\hat{\mathbf{a}}_2 J^n$$
$$\mathbf{a}_3^n = 4\hat{\mathbf{a}}_3 J^n$$

with

$$J = \frac{(\lambda^2 - \mu^2)(1 - \mu^2 \lambda^2)}{(\lambda^2 - \gamma^2)(1 - \gamma^2 \lambda^2)}, \ \mu \neq \{\pm \gamma, \pm \gamma^{-1}\}$$

it can be shown that $\mathfrak{U}_\lambda^{\mathbb{D}_2}=\bigoplus_k\{\mathbf{a}_1^k,\mathbf{a}_2^k,\mathbf{a}_3^k\}.$ Moreover

$$\begin{split} [\mathbf{a}_{1}^{n}, \mathbf{a}_{2}^{m}] &= \mathbf{a}_{3}^{n+m+1} + a(\gamma, \mu)\mathbf{a}_{3}^{n+m} \\ [\mathbf{a}_{3}^{n}, \mathbf{a}_{1}^{m}] &= 2\mathbf{a}_{1}^{n+m+1} + b(\gamma, \mu)\mathbf{a}_{1}^{n+m} - c(\gamma, \mu)\mathbf{a}_{2}^{n+m} \\ [\mathbf{a}_{3}^{n}, \mathbf{a}_{2}^{m}] &= -2\mathbf{a}_{2}^{n+m+1} - b(\gamma, \mu)\mathbf{a}_{2}^{n+m} + c(\gamma, \mu)\mathbf{a}_{1}^{n+m}. \end{split}$$
giving that $[A^{n}, A^{m}] \subset A^{n+m} \bigoplus A^{n+m+1}$

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$\mathfrak{U}=\textit{sl}(2,\mathbb{C})$

• Automorphic Lie algebras corresponding to \mathbb{Z}_N where Γ is a degenerated orbit are grading isomorphic for all N. For each algebra generators \mathbf{a}_1 , \mathbf{a}_2 , \mathbf{a}_3 and automorphic function J can be chosen such that the commutation relations take the form:

$$[\mathbf{a}_1, \mathbf{a}_2] = J\mathbf{a}_3$$
, $[\mathbf{a}_1, \mathbf{a}_3] = -2\mathbf{a}_1$, $[\mathbf{a}_2, \mathbf{a}_3] = 2\mathbf{a}_2$

 Autmorphic Lie algebras corresponding to Z_N where Γ = Z_N(γ) is a generic orbit are grading isomorphic for all N and γ. In this case the commutation relations take the form:

$$[\mathbf{a}_1, \mathbf{a}_2] = J^2 \mathbf{a}_3 - J \mathbf{a}_3$$
, $[\mathbf{a}_1, \mathbf{a}_3] = -2\mathbf{a}_1$, $[\mathbf{a}_2, \mathbf{a}_3] = 2\mathbf{a}_2$

• The automorphic Lie algebras corresponding to $\mathbb{D}_N, \mathbb{T}, \mathbb{O}, \mathbb{I}$ where Γ is a degenerated orbit of either group are all grading isomorphic to \mathbb{D}_2 with orbit $\Gamma = \{\infty, 0\}$. In this case there exist generators such that

$$[\mathbf{a}_1, \mathbf{a}_2] = \mathbf{a}_3$$
, $[\mathbf{a}_1, \mathbf{a}_3] = 4\mathbf{a}_2 - 2J\mathbf{a}_1$, $[\mathbf{a}_2, \mathbf{a}_3] = -4\mathbf{a}_1 + 2\mathbf{a}_2$.

For the degenerated orbit $G(\infty)$ the system arising from the automorphic Lie algebra corresponding to \mathbb{T} , \mathbb{O} and \mathbb{I} with 3x3 representation are point equivalent to the following system:

$$k_{1,t} = k_{1,xx} + k_{2,x}^2 + k_{2,x} (e^{-k_1 - k_2} + e^{-\omega k_1 - \omega^2 k_2} + e^{-\omega^2 k_1 - \omega k_2})$$

$$k_{2,t} = -k_{2,xx} - k_{1,x}^2 + k_{1,x} (e^{-k_1 - k_2} + \omega e^{-\omega k_1 - \omega^2 k_2} + \omega^2 e^{-\omega^2 k_1 - \omega k_2}).$$

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Darboux Transformations

- A Darboux transformation for a Lax operator L is a Matrix M such that $MLM^{-1} = \tilde{L}$.
- M maps solutions of $L\Psi = 0$ to solutions $\tilde{\Psi}$ of \tilde{L} . In fact we can see that $\tilde{L}\tilde{\Psi} = MLM^{-1}(M\Psi) = ML\Psi = 0$.
- Given $L = \partial_x + X$ we wish to find M such that $ML = \tilde{L}M$ so essentially we must solve the system:

$$M_{x}+\tilde{X}M-MX=0.$$

- When solving this system for Lax operators that have a particular symmetry we assume *M* has the same symmetry.
- This mapping is compatible with the system $L\Psi = 0$ in the sense that $\widetilde{(\Psi_x)} = \tilde{\Psi}_x$, for

$$\widetilde{(\Psi_{ imes})} = -\widetilde{X}\widetilde{\Psi} = -\widetilde{X}\widetilde{\Psi} = -\widetilde{X}M\Psi$$

$$\tilde{\Psi}_x = (M\Psi)_x = M_x\Psi + M\Psi_x = M_x\Psi - MU\Psi.$$

Darboux Transformations $sl(2, \mathbb{C})$

 \mathbb{Z}_2 reduction group

$$g_s: L(\lambda) \to \mathbf{Q}_s L(-\lambda) \mathbf{Q}_s^{-1},$$

• Degenerated orbit:

$$L = \partial_{x} + \lambda^{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} + \lambda \begin{pmatrix} 0 & 2p \\ 2q & 0 \end{pmatrix}$$
$$M = \lambda^{2} \begin{pmatrix} f & 0 \\ 0 & 0 \end{pmatrix} + \lambda \begin{pmatrix} 0 & fp \\ f\tilde{q} & 0 \end{pmatrix} + \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$$

• Generic orbit:

$$L = D_x + \frac{1}{\lambda - 1} S(p, q) - \frac{1}{\lambda + 1} \mathbf{Q}_s S(p, q) \mathbf{Q}_s,$$
$$S(p, q) := \frac{1}{p - q} \begin{pmatrix} p + q & -2pq \\ 2 & -p - q \end{pmatrix}.$$

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$$M = \frac{1}{\lambda - 1} f \begin{pmatrix} \tilde{q} & -p \tilde{q} \\ 1 & -p \end{pmatrix} - \frac{1}{\lambda + 1} \mathbf{Q}_{s} f \begin{pmatrix} \tilde{q} & -p \tilde{q} \\ 1 & -p \end{pmatrix} \mathbf{Q}_{s}^{-1} + \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix},$$

• Dihedral group \mathbb{D}_2 Lax operator corresponding to the degenerated orbit can be written as

$$L = D_{x} + \lambda^{2} \mathbf{Q}_{s} + \lambda \begin{pmatrix} 0 & 2p \\ 2q & 0 \end{pmatrix} + \frac{1}{\lambda} \begin{pmatrix} 0 & 2q \\ 2p & 0 \end{pmatrix} - \frac{1}{\lambda^{2}} \mathbf{Q}_{s},$$

$$M = f \left(\lambda^{2} \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} + \lambda \begin{pmatrix} 0 & p \\ \tilde{q} & 0 \end{pmatrix} + u \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} + \frac{1}{\lambda} \begin{pmatrix} 0 & \tilde{q} \\ p & 0 \end{pmatrix} + \frac{1}{\lambda^{2}} \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} \right)$$

Discrete Systems from Darboux Transformations

- Suppose we have Darboux transformations M and N such that $MLM^{-1} = \tilde{L}$ and $NLN^{-1} = \hat{L}$.
- M maps a fundamental solution Ψ to a solution $\tilde{\Psi}$.
- N maps a fundamental solution Ψ to a solution $\hat{\Psi}$.
- Imposing that $\tilde{\hat{\Psi}} = \hat{\tilde{\Psi}}$ and interpreting and as directions on a Lattice we obtain a difference system between the entries of our Darboux matrices.
- $\hat{\Psi} = \tilde{\Psi}$ is equivalent to the condition that $\hat{M}N = \tilde{N}M$.
- The differential difference equations obtained when deriving the Darboux transformations act as symmetries of our discrete system:

$$\begin{aligned} \frac{d}{dx}(\hat{M}N - \tilde{N}M) &= \hat{M}_{x}N + \hat{M}N_{x} - \tilde{N}_{x}M - \tilde{N}M_{x} \\ &= -\tilde{\hat{U}}(\hat{M}N - \tilde{N}M) + (\hat{M}N - \tilde{N}M)U \\ &= 0. \end{aligned}$$

Tetrahedral Group

For $L = \partial_x + X$ and $S = \partial_t + T$ we assume L and S are invariant with respect to the following group of transformations:

$$g_{s}: \mathbf{a}(\lambda) \to \mathbf{Q}_{s}\mathbf{a}(\sigma_{s}^{-1}(\lambda))\mathbf{Q}_{s}^{-1}$$
$$g_{r}: \mathbf{a}(\lambda) \to \mathbf{Q}_{r}\mathbf{a}(\sigma_{r}^{-1}(\lambda))\mathbf{Q}_{r}^{-1}$$
$$\sigma_{s}(\lambda) = \omega\lambda , \sigma_{r}(\lambda) = \frac{\lambda+2}{\lambda-1}.$$

Where $\omega = e^{\frac{2\pi i}{3}}$

$$\mathbf{Q}_{s} = \begin{pmatrix} \omega & 0 & 0 \\ 0 & \omega^{2} & 0 \\ 0 & 0 & 1 \end{pmatrix}, \mathbf{Q}_{r} = \begin{pmatrix} -1 & 2 & 2 \\ 2 & -1 & 2 \\ 2 & 2 & -1 \end{pmatrix}.$$

This group is Isomorphic to the tetrahedral group. It can be readily verified that $g_s^3 = g_r^2 = (g_r g_s)^3 = id$.

Generators of Tetrahedral automorphic Lie Algebra



Let $L = \partial_x + u_1 \mathbf{a}_1 + u_2 \mathbf{a}_2 + u_3 \mathbf{a}_3$ with $u_1 u_2 u_3 = 1$. In this case we may find Darboux transformations of the form

$$M = f\mathbf{I} + \alpha(u_1u_2\tilde{u}_1\mathbf{a}_1 + u_2\tilde{u}_1\tilde{u}_2\mathbf{a}_2 + \mathbf{a}_3)$$

where \tilde{u}_i , f and α satisfy particular differential relations, ie

$$\begin{aligned} \alpha' &= A(u_i, \tilde{u}_i, f) \\ f' &= B(u_i, \tilde{u}_i, f) \\ \tilde{u}_1' &= C(u_i, \tilde{u}_i, f) \\ \tilde{u}_2' &= D(u_i, \tilde{u}_i, f). \end{aligned}$$

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Degenerated cases

From direct computation we obtain that

$$\det[M] = EJ_1 + F_1 = EJ_2 + F_2$$

where

$$J_{1} = \frac{\lambda^{3}(\lambda^{3} + 8)^{3}}{4(\lambda^{3} - 1)^{3}}$$
$$J_{2} = \frac{(-8 - 20\lambda^{3} + \lambda^{6})^{2}}{4(\lambda^{3} - 1)^{3}}$$
$$E = \frac{1}{16}\alpha^{3}u_{1}u_{2}^{2}\tilde{u}_{1}^{2}\tilde{u}_{2}$$

 $F_{1} = \frac{1}{27} (3f + \alpha (1 + u_{1}u_{2}\tilde{u}_{1} - 2u_{2}\tilde{u}_{1}\tilde{u}_{2}))(3f + \alpha (1 - 2u_{1}u_{2}\tilde{u}_{1} + u_{2}\tilde{u}_{1}\tilde{u}_{2}))$ $(3f + \alpha (-2 + u_{1}u_{2}\tilde{u}_{1} + u_{2}\tilde{u}_{1}\tilde{u}_{2}))$ $F_{2} = \frac{1}{27} (3f + \alpha (1 + u_{1}u_{2}\tilde{u}_{1} + u_{2}\tilde{u}_{1}\tilde{u}_{2}))q(f, u_{1}, u_{2}, \tilde{u}_{1}, \tilde{u}_{2}),$

Continued..

Setting $u_1 = u$ and $u_2 = v$ we obtain the Darboux transformations:

$$M_i(u, v, u_i, v_i) = \beta_i(g_i \mathbf{I} + uvu_i \mathbf{a}_1 + vu_i v_i \mathbf{a}_2 + \mathbf{a}_3).$$

for i = 1, ..., 4 where u_i , v_i are \tilde{u} and \tilde{v} associated to $M_i \alpha^i g_i = f_i$ are the roots of F_1 or F_2 :

$$g_{1} = -\frac{1}{3}(-2 + uvu_{1} + vu_{1}v_{1})$$

$$g_{2} = -\frac{1}{3}(1 + uvu_{2} - 2vu_{2}v_{2})$$

$$g_{3} = \frac{1}{3}(-1 + 2uvu_{3} - vu_{3}v_{3})$$

$$g_{4} = -\frac{1}{3}(1 + uvu_{4} + vu_{4}v_{4}).$$

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Imposing the compatibility equations,

 $M_i(u_j, v_j, u_{ij}, v_{ij})M_j(u, v, u_j, v_j) - M_j(u_i, v_i, u_{ij}, v_{ij})M_i(u, v, u_i, v_i) = 0,$ we obtain the systems:

• For i=1, j=4

$$u_{4}v_{4}\alpha_{4}^{1}\alpha^{4} - u_{1}v_{1}\alpha^{1}\alpha_{1}^{4} = 0$$

$$u_{14}v_{4} + uv(-1 + u_{4}u_{14}v_{4} - u_{1}u_{14}v_{4}) = 0$$

$$uu_{1}^{2}vv_{4} + uv_{14} - u_{1}(v_{4}(1 + uv(u_{4} - v_{14})) + uvv_{1}v_{14}) = 0.$$
(1)
• For i=2, j=4

$$(-1 + u_{24}v_{4}v_{24})\alpha_{4}^{2}\alpha^{4} + (1 - u_{2}vv_{2})\alpha^{2}\alpha_{2}^{4} = 0$$

$$u_{4}v_{4}(1 - u_{2}vv_{2}) + u_{2}v_{2}(-1 + u_{24}v_{4}v_{24}) = 0$$

$$-u_{2}v_{2} + u(-u_{4} + u_{2} + v_{24}) = 0.$$
(2)

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Continued...

• For i=3, j=4

$$(-1 + u_4 u_{34} v_4) \alpha_4^3 \alpha^4 + (1 - u u_3 v) \alpha^3 \alpha_3^4 = 0$$

- $u_3 v_3 + u_4 v_4 (1 - u u_3 v + u_3 u_{34} v_3) = 0$ (3)
- $u_4 v_4 + (u + v_4 - v_3) v_{34} = 0.$

$$u_{1}v_{1}\alpha^{1}\alpha_{1}^{2} - u_{2}v_{2}\alpha_{2}^{1}\alpha^{2} = 0$$

- $u_{1}v + u_{12}v_{12} = 0$
- $uu_{1}^{2}vv_{2} - uv_{12} + u_{1}v_{2}(1 - u_{2}vv_{2} + uv(u_{2} + v_{12})) = 0.$ (4)

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$$- u_{3}v_{3}\alpha_{3}^{2}\alpha^{3} + u_{2}v_{2}\alpha^{2}\alpha_{2}^{3} = 0$$

$$- u_{2}v_{2} + uv_{23} = 0$$

$$- u^{2}u_{3}u_{2}vv_{2} - u_{2}^{2}u_{23}v_{3}v_{2}^{2} +$$

$$u(u_{2}v_{2} + u_{3}v_{3}(-1 + u_{2}(u_{23} + v)v_{2})) = 0.$$
(5)

$$- u_{1}v_{1}\alpha^{1}\alpha_{1}^{3} + u_{3}v_{3}\alpha_{3}^{1}\alpha^{3} = 0$$

$$- uv + u_{13}v_{3} = 0$$

$$- uu_{1}^{2}vv_{1} - uv_{13} + u_{1}(uv(u + v_{1})v_{13} + v_{3}(1 - uvv_{13})) = 0.$$
(6)

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Corresponding differential difference relations

$$-u + u_1 + u^2 u_1 v - u u_1^2 v + u_1 v u' + u v u'_1 + u u_1 v' = 0$$

$$-u v + 2u v_1 - u_1 v_1 - u^2 u_1 v v_1 + u u_1^2 v v_1 + u u_1 v'_1 = 0$$

$$uv - u_2v_2 + u^2u_2vv_2 - uu_2^2vv_2 - uu_2v^2v_2 - uu_2vv_2^2 + 2u_2^2vv_2^2 + u_2vv_2u' + uvv_2u'_2 + uu_2v_2v' = 0 - u^2v_2 + uu_2v_2 + uv_2^2 - u_2v_2^2 - v_2u' + uv_2' = 0$$

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• For i=3

$$uv - u^2 u_3 v^2 - u_3 v_3 + u u_3^2 v v_3 + u_3 v v_3 u' + u v v_3 u_3' + u u_3 v_3 v' = 0$$

- $u^2 v + u^2 v_3 + u v v_3 - u v_3^2 - v_3 u' + u v_3' = 0$

$$uv - u_4v_4 + u^2u_4vv_4 - uu_4^2vv_4 + u_4vv_4u' + uvv_4u'_4 + uu_4v_4v' = 0$$

- $u^2v_4 + uu_4v_4 + uvv_4 - uv_4^2 - v_4u' + uv'_4 = 0$

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- We can find Integrable pdes by considering reductions of general Lax pairs.
- We can impose a symmetry upon our Lax pairs to achieve this aim.
- Lax Pairs posessing such symmetries form a Lie algebra.
- For certain spectral dependence and reductions these Lie algebras are grading Isomorphic.
- Darboux transformations can be used to construct discrete systems.
- These discrete systems automatically possess a symmetry.