# Trajectories of the Plate-Ball Problem

# Vladimir Pulov<sup>1</sup> <u>Mariana Hadzhilazova</u>,<sup>2</sup> Ivailo Mladenov<sup>2</sup>

<sup>1</sup>Department of Physics, Technical University of Varna <sup>2</sup>Institute of Biophysics, Bulgarian Academy of Science

# Geometry, Integrability and Quantization June 3-8, 2016

イロト イヨト イヨト イヨト

### Trajectories of the Plate-Ball Problem

## 1. The Plate-Ball Problem

Statement of the Problem Mathematical Prerequisites Mathematical Model

### 2. Explicit Parametrization

Via the Jacobian Elliptic Functions and Elliptic Integrals The Plate-Ball Problem and the Eulerian Elastica Trajectories via Mathematica<sup>®</sup>

< 回 > < 三 > < 三 >

The Problem of a Rolling Sphere (David Kendall, Oxford, 1950s)

- A spherical ball rests on an infinite horizontal table.
- The ball has to be transferred from a given initial state to an arbitrary final state, meaning its position and orientation in space via a sequence of moves.
- Each move consists of rolling the ball along some straight line on the table.

(4月) (1日) (日)

• The axis of rotation must be horizontal and there must be no slipping between the ball and the table.

# The Problem of a Rolling Sphere (David Kendall, Oxford, 1950s)

How many moves, N, will be necessary and sufficient to reach any final state?

Ans: N = 3 (Hammersley, 1983)

< 回 > < 三 > < 三 >

# Rolling Along a Curve of Shortest Length (Hammersley, 1983)

▲御▶ ▲唐▶ ★唐▶

Which is the shortest curved path between the prescribed initial and final states?

## Quaternions and Rotations

A quaternion, q, is a combination of a scalar and a vector

 $\mathbf{q} = \rho_0 + \rho_1 \mathbf{i} + \rho_2 \mathbf{j} + \rho_3 \mathbf{k}$ 

A quaternion product,  ${\bf v}\,{\bf w},$  of two vectors  ${\bf v}$  and  ${\bf w}$  is a quaternion

 $\mathbf{v} \, \mathbf{w} = - \, \mathbf{v} \cdot \mathbf{w} + \mathbf{v} \times \mathbf{w}$ 

A rotation through angle  $\varphi$  about the unit vector **u** is represented by the quaternion pair

$$\pm q = \pm \mathrm{e}^{rac{1}{2}arphi \mathrm{u}} = \pm (\cos rac{1}{2} arphi + \mathrm{u} \sin rac{1}{2} arphi) = \pm \mathrm{u}_2 \mathrm{u}_1$$

A (2) > (

### **Optimal Control Problem**

State equation:  $\dot{\mathbf{x}} = f(\mathbf{x}, \mathbf{h}), \ \mathbf{x} \in \mathbb{R}^n, \ \mathbf{h} \in \mathbb{R}^r$ Control function:  $\mathbf{h} = \mathbf{h}(t) \in \Delta$  (admissible controls) Trajectory of the controlled system:  $\mathbf{x} = \mathbf{x}(t)$ Cost function:  $\mathcal{J}(\mathbf{h}) = \int_0^T f_0(t, \mathbf{x}(t), \mathbf{h}(t)) dt$ 

Given an initial point  $\mathbf{x}(0) = \mathbf{x}_0$  and a final point  $\mathbf{x}(T) = \mathbf{x}_E$ , find optimal control  $\mathbf{h}^*(t)$ , such that minimizes the cost function

 $\mathcal{J}(h^\star) \leq \mathcal{J}(h)$ 

< 回 > < 三 > < 三 >

Pontryagin Maximum Principle (free-final-time problem)

Optimal control  $\mathbf{h}^{\star}(t)$  and optimal trajectory  $\mathbf{x}^{\star}(t)$  are known. Hence, there exists adjoint function  $\lambda^{\star}(t) \in \mathbb{R}^{n+1}$  satisfying

$$\dot{\lambda} = -rac{\partial \mathcal{H}}{\partial \mathsf{x}}(\lambda(t),\mathsf{x}^{\star}(t),\mathsf{h}^{\star}(t))$$

where the Hamiltonian is given by

$$\mathcal{H}(\lambda,\mathbf{x},\mathbf{h}) = \sum_{\alpha=0}^{n} \lambda_{\alpha} f_{\alpha}(\mathbf{x},\mathbf{h})$$

and the Pontryagin maximum principle holds

$$\max_{h\in\Delta}\mathcal{H}(\lambda,\mathbf{x}^{\star},\mathbf{h})=\mathcal{H}(\lambda,\mathbf{x}^{\star},\mathbf{h}^{\star})\equiv0, \ \lambda_{0}\leq0$$

<日本<br />
<日本</p>

# Differential Equation of the Rolling Sphere (Hammersley, 1983)

The sphere is rolled along a unit speed curve  $\Gamma$ , parameterized by the arc length t. The resultant rotation at t is given by the quaternion

 $q(t) = 
ho_0(t) + 
ho_1(t)\mathbf{i} + 
ho_2(t)\mathbf{j} + 
ho_3(t)\mathbf{k}$ 

The instantaneous axis of rotation is specified by the unit vector h(t) along the axis of rotation

 $\mathbf{h}(t) = h_1(t)\mathbf{i} + h_2(t)\mathbf{j}$ 

The differential equation of the rolling sphere is of the form:

$$\dot{q} = \frac{1}{2}\mathbf{h} q$$

< 回 > < 三 > < 三 >

Rolling a Sphere Along a Curve of the Shortest Length (a "minimum-time" optimal control problem)

Cost function:  $\mathcal{J} \equiv T = \int_0^T 1.dt$  (*T* - length of the curve) Control function:  $\mathbf{h} = \mathbf{h}(t) \in \Delta$  ( $\Delta = \{\mathbf{h} : |\mathbf{h}| = 1\}$ ) Initial position and orientation: O(0,0), q(0) = (1,0,0,0)Final position:

$$x(T) \equiv \int_0^T h_2(t) \, \mathrm{d}t = x_E, \qquad z(T) \equiv \int_0^T h_1(t) \, \mathrm{d}t = -z_E$$

・ 同 ト ・ ヨ ト ・ ヨ ト …

Final orientation:  $q(T) = (\alpha_0, \alpha_1, \alpha_2, \alpha_3)$ 

## The Plate-Ball Problem Mathematical Model



æ

# The Plate-Ball Problem Mathematical Model

### State Variables and State Equations

State variables:  $\rho = (\rho_0, \rho_1, \rho_2, \rho_3, \rho_{-1}, \rho_4, \rho_5)$ 

 $(\rho_0, \rho_1, \rho_2, \rho_3)$  - coordinates of the quaternion q $(\rho_{-1}, \rho_4, \rho_5)$  - state variable inferred from the integrals  $\int_0^T 1 dt = T, \quad \int_0^T h_2(t) dt = x_E, \quad \int_0^T h_1(t) dt = -z_E$ 

State equations:

$$\dot{\rho}_0 = -\frac{1}{2}(h_1\rho_1 + h_2\rho_2), \qquad \dot{\rho}_1 = \frac{1}{2}(h_1\rho_0 + h_2\rho_3)$$
$$\dot{\rho}_2 = \frac{1}{2}(-h_1\rho_3 + h_2\rho_0), \qquad \dot{\rho}_3 = \frac{1}{2}(h_1\rho_2 - h_2\rho_1)$$
$$\dot{\rho}_{-1} = 1, \qquad \dot{\rho}_4 = \frac{1}{2}h_1, \qquad \dot{\rho}_5 = \frac{1}{2}h_2$$

・ 同 ト ・ ヨ ト ・ ヨ ト …

Adjoint Variables and the Hamiltonian (Arthurs and Walsh, 1986)

Adjoint variables:  $\lambda = (\lambda_{-1}, \lambda_0, \lambda_1, \lambda_2, \lambda_3, \lambda_4, \lambda_5)$ The Hamiltonian:

 $\mathcal{H} = \lambda_{-1} - \frac{1}{2}\lambda_0(h_1\rho_1 + h_2\rho_2) + \frac{1}{2}\lambda_1(h_1\rho_0 + h_2\rho_3) \\ + \frac{1}{2}\lambda_2(-h_1\rho_3 + h_2\rho_0) + \frac{1}{2}\lambda_3(h_1\rho_2 - h_2\rho_1) \\ + \frac{1}{2}\lambda_4h_1 + \frac{1}{2}\lambda_5h_2$ 

(日本) (日本) (日本)

Adjoint Equations (Arthurs and Walsh, 1986)

Vector form: 
$$\dot{\lambda} = -\frac{\partial \mathcal{H}}{\partial \rho}$$

Coordinate form:

$$\begin{split} \lambda_{-1} &= 0, & \lambda_4 = 0, & \lambda_5 = 0 \\ \dot{\lambda}_0 &= -\frac{1}{2}(h_1\lambda_1 + h_2\lambda_2), & \dot{\lambda}_1 = \frac{1}{2}(h_1\lambda_0 + h_2\lambda_3) \\ \dot{\lambda}_2 &= \frac{1}{2}(-h_1\lambda_3 + h_2\lambda_0), & \dot{\lambda}_3 = \frac{1}{2}(h_1\lambda_2 - h_2\lambda_1) \end{split}$$

・ロト ・四ト ・ヨト ・ヨト

3

Rolling a Sphere Along a Curve of the Shortest Length (a "minimum-time" optimal control problem)

Given the initial position x(0) = z(0) = 0 and the final position

$$\mathbf{x}(T) \equiv \int_0^T h_2(t) \,\mathrm{d}t = \mathbf{x}_E, \qquad \mathbf{z}(T) \equiv \int_0^T h_1(t) \,\mathrm{d}t = -\mathbf{z}_E$$

the initial and final orientation of the sphere

 $q(0) = (1, 0, 0, 0), \qquad q(T) = (\alpha_0, \alpha_1, \alpha_2, \alpha_3)$ 

find the optimal control  $(h_1^*, h_2^*)$  (the optimal unit vector  $\mathbf{h}^*$  along the axis of rotation) such that minimizes the cost function  $\mathcal{J} \equiv \mathcal{T}$  (minimizes the length  $\mathcal{T}$  of the curve).

・ロッ ・雪ッ ・ヨッ

# Optimal Controls (applying Pontryagin maximum principle)

The optimal controls  $(h_1^{\star}, h_2^{\star})$  are found by choosing

 $\lambda_{-1} = -1$ 

(as for "minimum-time" problem) and maximizing the Hamiltonian  $\mathcal{H}$  amongst the admissible controls

 $h_1^2 + h_2^2 - 1 = 0$ 

The optimal controls are found to be (Hammersley, 1983)

$$h_1^{\star} = \lambda_4 - \lambda_0 \rho_1 + \lambda_1 \rho_0 - \lambda_2 \rho_3 + \lambda_3 \rho_2$$
  
$$h_2^{\star} = \lambda_5 - \lambda_0 \rho_2 + \lambda_1 \rho_3 + \lambda_2 \rho_0 - \lambda_3 \rho_1$$

イロト 不得下 イヨト イヨト

Curvature Equation I (Arthurs and Walsh, 1986)

The coordinates of **h** are given by

 $h_1 = \cos\psi, \qquad h_1 = \sin\psi \tag{1}$ 

where  $\psi$  is the angle between **h** and the coordinate axis Ox.

The curvature  $\kappa$  of a plane curve  $\Gamma$  is given by

$$\kappa = \dot{\psi}$$
 (2)

・ 同下 ・ ヨト ・ ヨト

Based on the state and adjoint equations  $(\dot{\lambda} = -\frac{\partial \mathcal{H}}{\partial \rho}, \dot{\rho} = \frac{\partial \mathcal{H}}{\partial \lambda})$ , and the expression for  $(h_1^{\star}, h_2^{\star})$ , it is deduced from (1) and (2) that

$$\dot{\kappa} = \lambda \cos{(\psi + \varepsilon)}, \qquad \lambda, \varepsilon = \text{const}$$

Curvature Equation II (Arthurs and Walsh, 1986)

On integrating the last equation

 $\dot{\kappa} = \lambda \cos{(\psi + \varepsilon)}, \qquad \lambda > 0, \ \varepsilon = ext{const}$ 

and making use of

 $\dot{x} = \sin{(\psi + \varepsilon)}, \quad \dot{z} = -\cos{(\psi + \varepsilon)}, \quad \varepsilon = \text{const}$ 

it is obtained that the curvature of  $\Gamma$  satisfies the equation

 $\kappa = -\lambda z - \mu, \qquad \lambda > 0, \ \mu = \text{const}$ 

which is equivalently written as

$$rac{z''}{(1+{z'}^2)^{3/2}}=-\lambda z-\mu$$
 (z'=dz/dx)

(四) (日) (日)

## Curvature Equation III

On integrating the last equation

$$rac{z''}{(1+{z'}^2)^{3/2}}=-\lambda z-\mu$$
  $(z'={
m d} z/{
m d} x)$ 

under the initial condition

z'=0 for  $z=\eta=$  maximum deflection ordinate

it is obtained the equation

$$(z')^2 = rac{1 - \left[1 + rac{\lambda}{2}(z^2 - \eta^2) + \mu(z - \eta)
ight]^2}{\left[1 + rac{\lambda}{2}(z^2 - \eta^2) + \mu(z - \eta)
ight]^2}, \quad \lambda > 0, \ \mu, \ \eta = ext{const}$$

## Curvature Equation IV The Intrinsic Equation of the Trajectory

On introducing the arc length parameter t, it follows from

$$\kappa = -\lambda z - \mu$$

$$\left(\frac{\mathrm{d}z}{\mathrm{d}x}\right)^2 = \frac{1 - \left[1 + \frac{\lambda}{2}(z^2 - \eta^2) + \mu(z - \eta)\right]^2}{\left[1 + \frac{\lambda}{2}(z^2 - \eta^2) + \mu(z - \eta)\right]^2}$$

the intrinsic equation of the trajectory

$$\frac{\mathrm{d}\kappa}{\mathrm{d}t} = \frac{1}{4}(\sigma^2 - \kappa^2)(\kappa^2 + 4\lambda - \sigma^2), \quad \sigma = \mu + \eta\lambda$$

<日本<br />
<日本</p>

#### The Plate-Ball Problem and the Eulerian Elastica

On substituting with

$$z = \zeta - \frac{\mu}{\lambda}, \qquad \sigma^2 = 2(1-\nu)\lambda$$

where  $\zeta$  is a new coordinate and  $\nu < 1$  is a new parameter, the equations of the plate-ball problem are reduced to the equations of the Eulerian elastica (Djondjorov, Hadzhilazova, Mladenov and Vassilev, 2008)

$$\frac{\mathrm{d}x}{\mathrm{d}t} = \frac{\lambda\zeta^2}{2} + \nu$$

$$\left(rac{\mathrm{d}\zeta}{\mathrm{d}t}
ight)^2 = -rac{\lambda^2\zeta^4}{4} - \lambda\nu\zeta^2 - \nu^2 + 1$$

## Explicit Parametrization Via the Jacobian Elliptic Functions and Elliptic Integrals

Trajectories of the Plate-Ball Problem  $u \in (-1, 1)$ 

$$\mathbf{x}(t) = rac{2}{\sqrt{\lambda}} E(\mathrm{am}(\sqrt{\lambda}t,k),k) - t, \qquad \mathbf{z}(t) = \mathbf{a} \mathrm{cn}(\sqrt{\lambda}t,k) - rac{\mu}{\lambda}$$

where

$$a = \sqrt{\frac{2(1-\nu)}{\lambda}}, \qquad k = \sqrt{\frac{1-\nu}{2}}, \qquad \nu = 1 - \frac{\sigma^2}{2\lambda}$$

 $\sigma = \mu + \eta \lambda$ 

E(u,k) incomplete elliptic integral of second order am(u,k) Jacobian amplitude function cn(u,k) Jacobian elliptic cosine function

# Trajectories of the Plate-Ball Problem u = -1

$$\begin{aligned} x(t) &= \frac{4\tanh\left(\sqrt{\lambda}t\right)}{\sqrt{\lambda}} - t, \qquad z(t) = \frac{4\operatorname{sech}(\sqrt{\lambda}t)}{\sqrt{\lambda}} - \frac{\mu}{\lambda} \\ \nu &= 1 - \frac{\sigma^2}{2\lambda}, \qquad \sigma = \mu + \eta\lambda \end{aligned}$$

・ロト ・回ト ・ヨト ・ヨト

æ

## Explicit Parametrization Via the Jacobian Elliptic Functions and Elliptic Integrals

Trajectories of the Plate-Ball Problem u < -1

$$x(t) = aE(am(\sqrt{\frac{\lambda(1-\nu)}{2}}t,k),k) + \nu t$$

$$z(t) = a dn(\sqrt{rac{\lambda(1-
u)}{2}t},k) - rac{\mu}{\lambda}$$

#### where

$$a = \sqrt{rac{2(1-
u)}{\lambda}}, \quad k = \sqrt{rac{2}{1-
u}}, \quad 
u = 1 - rac{\sigma^2}{2\lambda}, \quad \sigma = \mu + \eta\lambda$$

E(u, k) incomplete elliptic integral of second order am(u, k) Jacobian amplitude function dn(u, k) Jacobian elliptic delta function

Vladimir Pulov, Mariana Hadzhilazova, Ivailo Mladenov Trajectories of the Plate-Ball Problem

# Case I

# A Trajectory of the Plate-Ball Problem (Arc of the Eulerian Elastica for $\lambda = 4, \sigma = 2$ )



### <u>Case II</u>

A Trajectory of the Plate-Ball Problem (Arc of the Eulerian Elastica for  $\lambda = 4, \sigma = 2.83$  )



A (1) > A (2) > A

## <u>Case III</u>

A Trajectory of the Plate-Ball Problem (Arc of the Eulerian Elastica for  $\lambda = 4, \sigma = 3.35$  )



Case IV

A Trajectory of the Plate-Ball Problem (Arc of the Eulerian Elastica for  $\lambda = 4, \sigma = 3.63564$  )



• • = • •

#### <u>Case V</u>

A Trajectory of the Plate-Ball Problem (Arc of the Eulerian Elastica for  $\lambda = 4, \sigma = 3.9$  )



▲□ → ▲ 三 → ▲ 三

### Explicit Parametrization Trajectories via Mathematica®

#### <u>Case VI</u>

# A Trajectory of the Plate-Ball Problem (Arc of the Eulerian Elastica for $\lambda = 4, \sigma = 4$ )



A (1) > A (2) > A

Vladimir Pulov, Mariana Hadzhilazova, Ivailo Mladenov Trajectories of the Plate-Ball Problem

#### Case VII

A Trajectory of the Plate-Ball Problem (Arc of the Eulerian Elastica for  $\lambda = 4, \sigma = 4.2$  )



▲□ → ▲ 三 → ▲ 三

#### References

- Hammersley J. (2008) Oxford Commemoration Ball, In Probability, Statistics and Analysis, London Math. Soc. Lecture Note Series, Cambridge University Press, **79**, 112-142.
- Arthurs A. and Walsh G. (1986) On Hammersley's Minimum Problem for a Rolling Sphere, Math. Proc. Camb. Phil. Soc., 529-534.
- Jurdjevic V.(1993) The geometry of the plate-ball problem, Arch. Rat. Mech. Anal. **124** 305-328.
- Djondjorov P., Hadzhilazova M., Mladenov I. and Vassilev V. (2008) *Explicit Parametrization of Euler's Elastica*, Geometry Integrability & Quantization. **9** 175-186.
- Mladenov I. and Hadzhilazova M. (2013) The Many Faces of Elastica (in Bulgarian), Avangard Prima, Sofia.
- Abramowitz M. and Stegun I. (1972) Handbook of Mathematical Functions, New York, Dover.