

ON THE DECOMPOSITION  
OF THE OSCILLATOR REPRESENTATION

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Following Howe [7] we shall call the oscillator representation in  $\mathbf{R}^d$  the representation of the Lie algebra  $f\mathfrak{U}(\epsilon, \mathbf{R})$  on the Schwartz space  $S(\mathbf{R}^d)$  or the corresponding (i.e. exponentiated) unitary representation on  $L^2(\mathbf{R}^d)$  of the double cover  $\widetilde{SL}(2, \mathbf{R})$  of  $SL(2, \mathbf{R})$  obtained in the following way. Let

$$h = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}, \quad e^+ = \frac{1}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad e^- = \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix}$$

be a standard basis of the Lie algebra  $sl(2, \mathbf{R})$  satisfying the commutation relations

$$[h, e^\pm] = \pm 2e^\pm, \quad [e^+, e^-] = h.$$

We let  $\Delta = \sum_{j=1}^d \frac{\partial^2}{\partial x_j^2}$  stand for the Laplace operator and we shall use the familiar notation  $r^2 = r^2(x) = x \cdot x$  for the function on  $\mathbf{R}^d$  equal to the square of the radius (for the standard euclidean metric). If  $D$  is a differential operator (e.g. the Laplacian) and  $f$  a smooth function, we shall write  $D \circ f$  for the composition of  $D$  with the operator of multiplication by  $f$ . The oscillator representation  $\omega^d$  of the Lie algebra  $sl(2, \mathbf{R})$  is defined by setting

$$\begin{aligned} \omega^d(h) &= \sum_{j=1}^d x_j \frac{\partial}{\partial x_j} + \frac{d}{2}, \\ \omega^d(e^+) &= \frac{i}{2} \sum_{j=1}^d x_j^2 = \frac{i}{2} r^2, \\ \omega^d(e^-) &= \frac{i}{2} \sum_{j=1}^d \frac{\partial^2}{\partial x_j^2} = \frac{i}{2} \Delta \end{aligned} \tag{1}$$

and extending by linearity to the whole of  $sl(2, \mathbf{R})$ . The fact that this representation of the Lie algebra  $sl(2, \mathbf{R})$  is a derived representation of a unitary representation of the double cover  $\widetilde{SL}(2, \mathbf{R})$  of  $SL(2, \mathbf{R})$  is a part of a theorem of Shale and A. Weil asserting the existence of the metaplectic representation, cf. e.g. [6] or [3].

Many properties of the oscillator representation are easily accessed through the use of another basis satisfying the standard commutation relations, namely the one given by the matrices

$$k = \begin{pmatrix} 0 & i \\ -i & 0 \end{pmatrix}, \quad n^+ = \frac{1}{2} \begin{pmatrix} 1 & i \\ i & -1 \end{pmatrix}, \quad n^- = \frac{1}{2} \begin{pmatrix} 1 & -i \\ -i & -1 \end{pmatrix},$$

which satisfy

$$[k, n^\pm] = \pm 2n^\pm, \quad [n^+, n^-] = k.$$